

## EE2 Mathematics

### Solutions to Example Sheet 4: Fourier Transforms

1) Because  $f(t) = e^{-|t|} = \begin{cases} e^{-t}, & t > 0 \\ e^t, & t < 0 \end{cases}$  the Fourier transform of  $f(t)$  is

$$\bar{f}(\omega) = \int_{-\infty}^{\infty} e^{-i\omega t - |t|} dt = \int_{-\infty}^0 e^{t(1-i\omega)} dt + \int_0^{\infty} e^{-t(1+i\omega)} dt = \frac{2}{1+\omega^2}$$

2) (i) Designate  $\mathcal{F}\{f(t)\} = \bar{f}(\omega)$  with  $a$  a real constant of either sign. Then  $\mathcal{F}\{f(at)\} = \int_{-\infty}^{\infty} e^{-i\omega t} f(at) dt$ . Define  $\tau = at$  so  $d\tau = a dt$ . When  $a > 0$  the limits  $(-\infty, \infty)$  for  $\tau$  correspond to those for  $t$ , but when  $a < 0$  the direction reverses. Thus

$$\mathcal{F}\{f(at)\} = |a|^{-1} \int_{-\infty}^{\infty} e^{-i(\frac{\omega}{a})\tau} f(\tau) d\tau = |a|^{-1} \bar{f}\left(\frac{\omega}{a}\right)$$

(ii) The ‘shift property’ (in the formula sheets)  $\mathcal{F}\{f(t-a)\} = e^{-i\omega a} \bar{f}(\omega)$  can simply be proved by defining  $\tau = t-a$  then

$$\mathcal{F}\{f(t-a)\} = \int_{-\infty}^{\infty} e^{-i\omega t} f(t-a) dt = e^{-i\omega a} \int_{-\infty}^{\infty} e^{-i\omega \tau} f(\tau) d\tau = e^{-i\omega a} \bar{f}(\omega)$$

3) To find the Fourier transform of the non-normalized Gaussian  $f(t) = e^{-t^2}$  we first complete the square in the exponential

$$\bar{f}(\omega) = \int_{-\infty}^{\infty} e^{-i\omega t - t^2} dt = e^{-\frac{1}{4}\omega^2} \int_{-\infty}^{\infty} e^{-(t+\frac{1}{2}i\omega)^2} dt = \sqrt{\pi} e^{-\frac{1}{4}\omega^2}$$

The normalized auto-correlation function of  $e^{-t^2}$  is

$$\gamma(t) = \frac{\int_{-\infty}^{\infty} e^{-u^2} e^{-(t-u)^2} du}{\int_{-\infty}^{\infty} e^{-2u^2} du} = \frac{e^{-\frac{1}{2}t^2} \int_{-\infty}^{\infty} e^{-2(u-\frac{1}{2}t)^2} du}{\int_{-\infty}^{\infty} e^{-2u^2} du} = e^{-\frac{1}{2}t^2}$$

The integrals in the numerator & denominator cancel because they are equal; the origin of the former is shifted w.r.t. to the latter on the infinite  $u$ -axis but its value is not affected.

4) With  $f(t) = e^{-at^2}$  and  $g(t) = e^{-bt^2}$ , a minor re-scaling of the results of Q3 shows that

$$\bar{f}(\omega) = \sqrt{\frac{\pi}{a}} e^{-\frac{\omega^2}{4a}} \qquad \bar{g}(\omega) = \sqrt{\frac{\pi}{b}} e^{-\frac{\omega^2}{4b}}$$

The convolution theorem says that  $\mathcal{F}\left\{\int_{-\infty}^{\infty} f(t')g(t-t') dt'\right\} = \bar{f}(\omega)\bar{g}(\omega)$ . Therefore, with

$$h(t) = \int_{-\infty}^{\infty} \exp[-at'^2] \exp[-b(t-t')^2] dt' = f * g$$

we have  $\bar{h}(\omega) = \frac{\pi}{\sqrt{ab}} e^{-\frac{\omega^2(a+b)}{4ab}}$ .

5) With  $f(t) = e^{-t^2}$  for which  $\bar{f}(\omega) = \sqrt{\pi}e^{-\frac{1}{4}\omega^2}$  and  $g(t) = \cos at$  for which

$$\begin{aligned}\bar{g}(\omega) &= \int_{-\infty}^{\infty} e^{-i\omega t} \cos at \, dt = \frac{1}{2} \int_{-\infty}^{\infty} e^{-i\omega t} (e^{iat} + e^{-iat}) \, dt \\ &= \frac{1}{2} \int_{-\infty}^{\infty} \left\{ e^{-it(\omega-a)} + e^{-it(\omega+a)} \right\} \, dt = \pi \{ \delta(\omega - a) + \delta(\omega + a) \}\end{aligned}$$

Thus

$$\int_{-\infty}^{\infty} f(t)g(t) \, dt = \frac{\pi\sqrt{\pi}}{2\pi} \int_{-\infty}^{\infty} e^{-\frac{1}{4}\omega^2} \{ \delta(\omega - a) + \delta(\omega + a) \} \, d\omega = \sqrt{\pi} e^{-\frac{1}{4}a^2}.$$

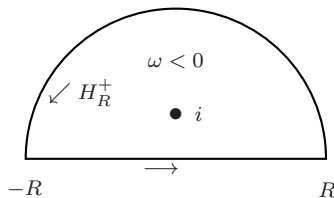
6) Write  $f(t) = (1 + t^2)^{-1}$  so that

$$\int_{-\infty}^{\infty} \frac{dt}{(1 + t^2)^2} = \int_{-\infty}^{\infty} |f(t)|^2 dt = \frac{1}{2\pi} \int_{-\infty}^{\infty} |\bar{f}(\omega)|^2 d\omega$$

Hence we want to evaluate  $\int_{-\infty}^{\infty} |\bar{f}(\omega)|^2 d\omega$ . To do this we must first find  $\bar{f}(\omega) = \int_{-\infty}^{\infty} \frac{e^{-i\omega t} dt}{1+t^2}$ . To apply Jordan's Lemma it is necessary to consider the two separate cases  $\omega < 0$  and  $\omega > 0$ .

(i)  $\omega < 0$ : Consider the complex integral  $\oint_{C_U} \frac{e^{-i\omega z} dz}{1+z^2}$  with  $C_U$  a semi-circle in the *upper*  $\frac{1}{2}$ -plane in which there is a simple pole at  $z = i$ . The residue at this pole is  $e^{\omega}/2i$  and the integral over  $H_R^+ \rightarrow 0$  as  $R \rightarrow \infty$  by Jordan's Lemma.

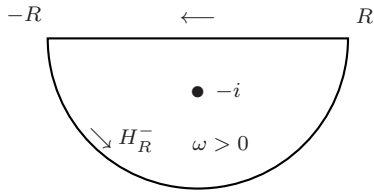
For the contour  $C_U$  in the upper  $\frac{1}{2}$ -plane ( $\omega < 0$ ):



$$\pi e^{\omega} = \oint_{C_U} \frac{e^{-i\omega z} dz}{1+z^2} = \int_{-R}^R \frac{e^{-i\omega x} dx}{1+x^2} + \int_{H_R^+} \frac{e^{-i\omega z} dz}{1+z^2}$$

(ii)  $\omega > 0$ : Consider the complex integral  $\oint_{C_L} \frac{e^{-i\omega z} dz}{1+z^2}$  with  $C_L$  a semi-circle in the *lower*  $\frac{1}{2}$ -plane in which there is a simple pole at  $z = -i$ . The residue at this pole is  $-e^{-\omega}/2i$  & the integral over  $H_R^- \rightarrow 0$  as  $R \rightarrow \infty$  by Jordan's Lemma.

For the contour  $C_L$  in the lower  $\frac{1}{2}$ -plane ( $\omega > 0$ ):



$$-\pi e^{-\omega} = \oint_{C_L} \frac{e^{-i\omega z} dz}{1+z^2} = \int_R^{-R} \frac{e^{-i\omega x} dx}{1+x^2} + \int_{H_R^-} \frac{e^{-i\omega z} dz}{1+z^2}$$

Note the reverse order of the limits in the real integral.

Thus, in the limit  $R \rightarrow \infty$ , we have  $\bar{f}(\omega) = \begin{cases} \pi e^{-\omega}, & \omega > 0 \\ \pi e^{\omega}, & \omega < 0 \end{cases}$ . Finally we can now calculate

$$\frac{1}{2\pi} \int_{-\infty}^{\infty} |\bar{f}(\omega)|^2 d\omega = \frac{\pi^2}{2\pi} \left\{ \int_{-\infty}^0 e^{2\omega} d\omega + \int_0^{\infty} e^{-2\omega} d\omega \right\} = \frac{1}{2}\pi.$$