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M4A32: Vortex dynamics  
Problem Sheet 4: SOLUTIONS  
(Point vortex dynamics on a sphere)

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**1.** On use of the facts that

$$\frac{\partial \psi}{\partial \theta} \Big|_{\phi} = \frac{\partial \psi}{\partial \zeta} \Big|_{\bar{\zeta}} \frac{\partial \zeta}{\partial \theta} \Big|_{\phi} + \frac{\partial \psi}{\partial \bar{\zeta}} \Big|_{\zeta} \frac{\partial \bar{\zeta}}{\partial \theta} \Big|_{\phi}, \quad \frac{\partial \psi}{\partial \phi} \Big|_{\theta} = \frac{\partial \psi}{\partial \zeta} \Big|_{\bar{\zeta}} \frac{\partial \zeta}{\partial \phi} \Big|_{\theta} + \frac{\partial \psi}{\partial \bar{\zeta}} \Big|_{\zeta} \frac{\partial \bar{\zeta}}{\partial \phi} \Big|_{\theta}, \quad (1)$$

and the relation

$$\zeta = \cot(\theta/2) e^{i\phi} \quad (2)$$

it follows that

$$\frac{\partial \zeta}{\partial \theta} \Big|_{\phi} = -\frac{\zeta}{\sin \theta}, \quad \frac{\partial \bar{\zeta}}{\partial \theta} \Big|_{\phi} = -\frac{\bar{\zeta}}{\sin \theta}, \quad \frac{\partial \zeta}{\partial \phi} \Big|_{\theta} = i\zeta, \quad \frac{\partial \bar{\zeta}}{\partial \phi} \Big|_{\theta} = -i\bar{\zeta}. \quad (3)$$

Now, on use of the facts that

$$\frac{\partial \psi}{\partial \theta} \Big|_{\phi} = \frac{\partial \psi}{\partial \zeta} \Big|_{\bar{\zeta}} \frac{\partial \zeta}{\partial \theta} \Big|_{\phi} + \frac{\partial \psi}{\partial \bar{\zeta}} \Big|_{\zeta} \frac{\partial \bar{\zeta}}{\partial \theta} \Big|_{\phi}, \quad \frac{\partial \psi}{\partial \phi} \Big|_{\theta} = \frac{\partial \psi}{\partial \zeta} \Big|_{\bar{\zeta}} \frac{\partial \zeta}{\partial \phi} \Big|_{\theta} + \frac{\partial \psi}{\partial \bar{\zeta}} \Big|_{\zeta} \frac{\partial \bar{\zeta}}{\partial \phi} \Big|_{\theta}, \quad (4)$$

then

$$\frac{\partial \psi}{\partial \theta} \Big|_{\phi} = -\frac{\zeta}{\sin \theta} \frac{\partial \psi}{\partial \zeta} - \frac{\bar{\zeta}}{\sin \theta} \frac{\partial \psi}{\partial \bar{\zeta}} \quad (5)$$

so that

$$\sin \theta \frac{\partial \psi}{\partial \theta} \Big|_{\phi} = -\zeta \frac{\partial \psi}{\partial \zeta} - \bar{\zeta} \frac{\partial \psi}{\partial \bar{\zeta}}. \quad (6)$$

It follows that

$$\begin{aligned} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial \psi}{\partial \theta} \right) &= -\frac{\zeta}{\sin \theta} \frac{\partial}{\partial \zeta} \left( -\zeta \frac{\partial \psi}{\partial \zeta} - \bar{\zeta} \frac{\partial \psi}{\partial \bar{\zeta}} \right) - \frac{\bar{\zeta}}{\sin \theta} \frac{\partial}{\partial \bar{\zeta}} \left( -\zeta \frac{\partial \psi}{\partial \zeta} - \bar{\zeta} \frac{\partial \psi}{\partial \bar{\zeta}} \right) \\ &= \frac{1}{\sin \theta} \left[ \zeta \frac{\partial \psi}{\partial \zeta} + \zeta^2 \frac{\partial^2 \psi}{\partial \zeta^2} + 2\zeta \bar{\zeta} \frac{\partial^2 \psi}{\partial \zeta \partial \bar{\zeta}} + \bar{\zeta}^2 \frac{\partial^2 \psi}{\partial \bar{\zeta}^2} + \bar{\zeta} \frac{\partial \psi}{\partial \bar{\zeta}} \right]. \end{aligned} \quad (7)$$

Also

$$\frac{\partial \psi}{\partial \phi} \Big|_{\theta} = i\zeta \frac{\partial \psi}{\partial \zeta} - i\bar{\zeta} \frac{\partial \psi}{\partial \bar{\zeta}}, \quad (8)$$

so that

$$\begin{aligned}\frac{\partial^2 \psi}{\partial \phi^2} &= i\zeta \frac{\partial}{\partial \zeta} \left( i\zeta \frac{\partial \psi}{\partial \zeta} - i\bar{\zeta} \frac{\partial \psi}{\partial \bar{\zeta}} \right) - i\bar{\zeta} \frac{\partial}{\partial \bar{\zeta}} \left( i\zeta \frac{\partial \psi}{\partial \zeta} - i\bar{\zeta} \frac{\partial \psi}{\partial \bar{\zeta}} \right) \\ &= -\zeta \frac{\partial \psi}{\partial \zeta} - \zeta^2 \frac{\partial^2 \psi}{\partial \zeta^2} + 2\zeta \bar{\zeta} \frac{\partial^2 \psi}{\partial \zeta \partial \bar{\zeta}} - \bar{\zeta}^2 \frac{\partial^2 \psi}{\partial \bar{\zeta}^2} - \bar{\zeta} \frac{\partial \psi}{\partial \bar{\zeta}}.\end{aligned}\quad (9)$$

By adding appropriate multiples of (7) and (9), we arrive at

$$\nabla_{\Sigma}^2 \psi = \left( \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2 \psi}{\partial \phi^2} \right) = \frac{4\zeta \bar{\zeta}}{\sin^2 \theta} \frac{\partial^2 \psi}{\partial \zeta \partial \bar{\zeta}}. \quad (10)$$

But

$$\sin^2 \theta = \frac{4\zeta \bar{\zeta}}{(1 + \zeta \bar{\zeta})^2} \quad (11)$$

so

$$\nabla_{\Sigma}^2 \psi = (1 + \zeta \bar{\zeta})^2 \frac{\partial^2 \psi}{\partial \zeta \partial \bar{\zeta}}. \quad (12)$$

## 2. In spherical polar coordinates

$$\nabla \cdot \mathbf{u} = \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} (u_{\theta} \sin \theta) + \frac{1}{\sin \theta} \frac{\partial u_{\phi}}{\partial \phi} = 0. \quad (13)$$

So we define

$$u_{\phi} = -\frac{\partial \psi}{\partial \theta}, \quad u_{\theta} = \frac{1}{\sin \theta} \frac{\partial \psi}{\partial \phi}. \quad (14)$$

Therefore

$$u_{\phi} - iu_{\theta} = -\frac{\partial \psi}{\partial \theta} - \frac{i}{\sin \theta} \frac{\partial \psi}{\partial \phi} \quad (15)$$

But we know, from Q1, that

$$\begin{aligned}\frac{\partial \psi}{\partial \theta} &= -\frac{\zeta}{\sin \theta} \frac{\partial \psi}{\partial \zeta} - \frac{\bar{\zeta}}{\sin \theta} \frac{\partial \psi}{\partial \bar{\zeta}}, \\ \frac{\partial \psi}{\partial \phi} &= i\zeta \frac{\partial \psi}{\partial \zeta} - i\bar{\zeta} \frac{\partial \psi}{\partial \bar{\zeta}},\end{aligned}\quad (16)$$

Therefore

$$u_{\phi} - iu_{\theta} = \frac{\zeta}{\sin \theta} \frac{\partial \psi}{\partial \zeta} + \frac{\bar{\zeta}}{\sin \theta} \frac{\partial \psi}{\partial \bar{\zeta}} + \frac{\zeta}{\sin \theta} \frac{\partial \psi}{\partial \zeta} - \frac{\bar{\zeta}}{\sin \theta} \frac{\partial \psi}{\partial \bar{\zeta}} = \frac{2\zeta}{\sin \theta} \frac{\partial \psi}{\partial \zeta}. \quad (17)$$

3. Consider

$$\zeta = \cot(\theta/2)e^{i\phi}, \quad \zeta' = \cot(\theta'/2)e^{i\phi'}. \quad (18)$$

and

$$\frac{2(\zeta - \zeta')(\bar{\zeta} - \bar{\zeta}')}{(1 + \zeta\bar{\zeta})(1 + \zeta'\bar{\zeta}')}.$$
(19)

Now

$$1 + \zeta\bar{\zeta} = 1 + \cot^2(\theta/2) = \operatorname{cosec}^2(\theta/2), \quad 1 + \zeta'\bar{\zeta}' = 1 + \cot^2(\theta'/2) = \operatorname{cosec}^2(\theta'/2)$$
(20)

while

$$\begin{aligned} (\zeta - \zeta')(\bar{\zeta} - \bar{\zeta}') &= (\cot(\theta/2)e^{i\phi} - \cot(\theta'/2)e^{i\phi'})(\cot(\theta/2)e^{-i\phi} - \cot(\theta'/2)e^{-i\phi'}) \\ &= \cot^2(\theta/2) + \cot^2(\theta'/2) - \cot(\theta/2)\cot(\theta'/2) \left[ e^{i(\phi'-\phi)} + e^{-i(\phi'-\phi)} \right] \end{aligned}$$
(21)

Therefore

$$\begin{aligned} \frac{2(\zeta - \zeta')(\bar{\zeta} - \bar{\zeta}')}{(1 + \zeta\bar{\zeta})(1 + \zeta'\bar{\zeta}')} &= 2 \sin^2(\theta/2) \sin^2(\theta'/2) \left[ \frac{\cos^2(\theta/2)}{\sin^2(\theta/2)} + \frac{\cos^2(\theta'/2)}{\sin^2(\theta'/2)} \right. \\ &\quad \left. - \frac{2 \cos(\theta/2) \cos(\theta'/2)}{\sin(\theta/2) \sin(\theta'/2)} \cos(\phi - \phi') \right]. \end{aligned}$$
(22)

This can be written

$$\begin{aligned} \frac{2(\zeta - \zeta')(\bar{\zeta} - \bar{\zeta}')}{(1 + \zeta\bar{\zeta})(1 + \zeta'\bar{\zeta}')} &= 2 \left[ \cos^2(\theta/2) \sin^2(\theta'/2) + \cos^2(\theta'/2) \sin^2(\theta/2) \right. \\ &\quad \left. - 2[\sin(\theta/2) \cos(\theta/2)][\sin(\theta'/2) \cos(\theta'/2)] \cos(\phi - \phi') \right]. \end{aligned}$$
(23)

But, from the half-angle formulae,

$$\begin{aligned} 1 - \cos \theta \cos \theta' &= (\cos^2(\theta/2) + \sin^2(\theta/2))(\cos^2(\theta'/2) + \sin^2(\theta'/2)) \\ &\quad - (\cos^2(\theta/2) - \sin^2(\theta/2))(\cos^2(\theta'/2) - \sin^2(\theta'/2)) \quad (24) \\ &= 2 [\cos^2(\theta'/2) \sin^2(\theta/2) + \cos^2(\theta/2) \sin^2(\theta'/2)]. \end{aligned}$$

By substituting for this in (23) we arrive at

$$\frac{2(\zeta - \zeta')(\bar{\zeta} - \bar{\zeta}')}{(1 + \zeta\bar{\zeta})(1 + \zeta'\bar{\zeta}')} = 1 - \cos \theta \cos \theta' - \sin \theta \sin \theta' \cos(\phi - \phi') = 1 - \cos \gamma. \quad (25)$$

**4.** Suppose there are  $N$  equal strength vortices at positions  $a, a\omega, a\omega^2, \dots, a\omega^{N-1}$  where

$$\omega^N = 1. \quad (26)$$

The streamfunction in a frame of reference corotating with some angular velocity  $\Omega$  is

$$\psi = -\frac{\kappa}{2} \log \left[ \frac{(\zeta^N - a^N)(\bar{\zeta}^N - a^N)}{(1 + \zeta\bar{\zeta})^N (1 + a^2)^N} \right] + \frac{2\Omega}{(1 + \zeta\bar{\zeta})}. \quad (27)$$

The velocity field is

$$u_\phi - iu_\theta = \frac{2\zeta}{\sin \theta} \left[ -\frac{N\kappa\zeta^{N-1}}{2(\zeta^N - a^N)} + \frac{N\kappa\bar{\zeta}}{2(1 + \zeta\bar{\zeta})} - \frac{2\Omega\bar{\zeta}}{(1 + \zeta\bar{\zeta})^2} \right] \quad (28)$$

For the configuration to be steady we must subtract off the velocity due to the vortex at  $\zeta = a$  and set the remaining terms equal to zero. The velocity due to the point vortex at  $\zeta = a$  is

$$\frac{2\zeta}{\sin \theta} \left( -\frac{\kappa}{2} \left[ \frac{1}{\zeta - a} - \frac{\bar{\zeta}}{1 + \zeta\bar{\zeta}} \right] \right). \quad (29)$$

so the condition for steadiness reduces to

$$\left[ -\frac{N\kappa\zeta^{N-1}}{2(\zeta^N - a^N)} + \frac{\kappa}{2(\zeta - a)} + \frac{(N-1)\kappa\bar{\zeta}}{2(1 + \zeta\bar{\zeta})} - \frac{2\Omega\bar{\zeta}}{(1 + \zeta\bar{\zeta})^2} \right] \Big|_{\zeta=a} = 0. \quad (30)$$

Using the facts that

$$\begin{aligned} \zeta^N &= a^N + Na^{N-1}(\zeta - a) + \frac{N(N-1)a^{N-2}}{2!}(\zeta - a)^2 + \dots \\ \zeta^{N-1} &= a^{N-1} + (N-1)a^{N-2}(\zeta - a) + \dots \end{aligned} \quad (31)$$

the condition (30) becomes

$$-\frac{\kappa(N-1)}{4a} + \frac{\kappa a(N-1)}{2(1+a^2)} - \frac{2\Omega a}{(1+a^2)^2} = 0. \quad (32)$$

On rearrangement, we deduce that

$$\Omega = \frac{\kappa(N-1)(a^4-1)}{8a^2}. \quad (33)$$

This can also be written as

$$\Omega = \frac{\kappa(N-1)\cos\theta_a}{2\sin^2\theta_a} \quad (34)$$

where  $a = \cot(\theta_a/2)$ .

If the vortices are on the equator then  $\theta_a = \pi/2$  and  $\Omega = 0$  so the configuration is stationary.

**5.** If fluid is in solid body rotation about an axis through the north pole with angular velocity  $\Omega$  then

$$u_\phi = \Omega \sin \theta, \quad u_\theta = 0 \quad (35)$$

so that

$$u_\phi - iu_\theta = \Omega \sin \theta = \frac{2\zeta}{\sin \theta} \frac{\partial \psi}{\partial \zeta}. \quad (36)$$

This means that

$$\frac{\partial \psi}{\partial \zeta} = \frac{2\Omega \bar{\zeta}}{(1+\zeta \bar{\zeta})^2}. \quad (37)$$

But

$$\omega = -\nabla_{\Sigma}^2 \psi = -(1+\zeta \bar{\zeta})^2 \frac{\partial^2 \psi}{\partial \zeta \partial \bar{\zeta}}, \quad (38)$$

so since

$$\frac{\partial^2 \psi}{\partial \zeta \partial \bar{\zeta}} = \frac{2\Omega}{(1+\zeta \bar{\zeta})^2} - \frac{4\Omega \zeta \bar{\zeta}}{(1+\zeta \bar{\zeta})^3} = \frac{2\Omega(1-\zeta \bar{\zeta})}{(1+\zeta \bar{\zeta})^3} \quad (39)$$

then

$$\omega = -(1+\zeta \bar{\zeta})^2 \frac{\partial^2 \psi}{\partial \zeta \partial \bar{\zeta}} = -\frac{2\Omega(1-\zeta \bar{\zeta})}{(1+\zeta \bar{\zeta})} = 2\Omega \cos \theta. \quad (40)$$

6. The required streamfunction is

$$\psi(\zeta, \bar{\zeta}) = \begin{cases} -\omega_S \log(1 + \zeta \bar{\zeta}), & |\zeta| < 1, \\ -\omega_N \log[(1 + \zeta \bar{\zeta})/(\zeta \bar{\zeta})], & |\zeta| > 1. \end{cases}$$

To verify this, on differentiation we find

$$\frac{\partial \psi}{\partial \zeta} = \begin{cases} -\omega_S \bar{\zeta}/(1 + \zeta \bar{\zeta}), & |\zeta| < 1, \\ -\omega_N \bar{\zeta}/(1 + \zeta \bar{\zeta}) + \omega_N/\zeta, & |\zeta| > 1, \end{cases}$$

and

$$\frac{\partial^2 \psi}{\partial \zeta \partial \bar{\zeta}} = \begin{cases} -\omega_S/(1 + \zeta \bar{\zeta})^2, & |\zeta| < 1, \\ -\omega_N/(1 + \zeta \bar{\zeta})^2, & |\zeta| > 1, \end{cases}$$

so that

$$\nabla_{\Sigma}^2 \psi = (1 + \zeta \bar{\zeta})^2 \frac{\partial^2 \psi}{\partial \zeta \partial \bar{\zeta}} = \begin{cases} -\omega_S, & |\zeta| < 1, \\ -\omega_N, & |\zeta| > 1. \end{cases}$$

Continuity of velocity at  $|\zeta| = 1$  requires continuity of  $\partial \psi / \partial \zeta$  there which implies

$$-\omega_S \bar{\zeta}/(1 + \zeta \bar{\zeta}) = -\omega_N \bar{\zeta}/(1 + \zeta \bar{\zeta}) + \omega_N/\zeta$$

which, on the equator  $|\zeta| = 1$ , implies

$$\omega_N = -\omega_S.$$

This is consistent with the Gauss constraint that the total vorticity over the sphere is zero.