

Classical Density Functional Theory for Intergranular Films

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Department of Physics
Imperial College London

INCEMS Meeting, Stuttgart, March 2006

Outline

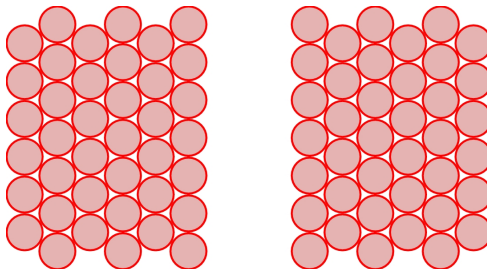
1 Motivation – Model, Observables, Parameters

- Model
- Observables and Parameters

2 Classical Density Functional Theory

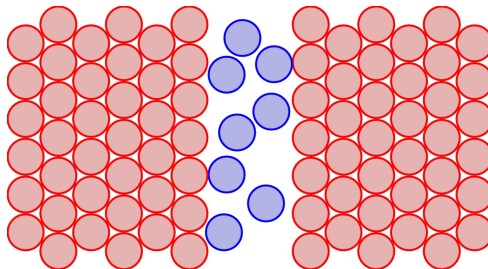
- Principles
- An Example
- From DFT to Phase Field Models

Model



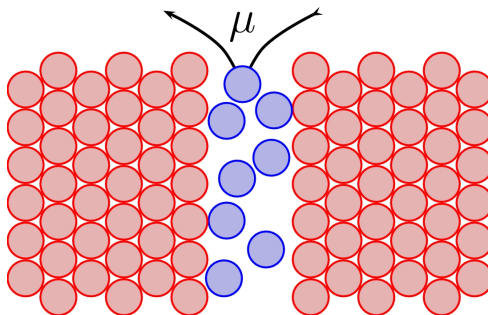
- Two **confining** lattices (3D)
- Interface: **liquid layer** in between
- Reservoir: **chemical potential** μ

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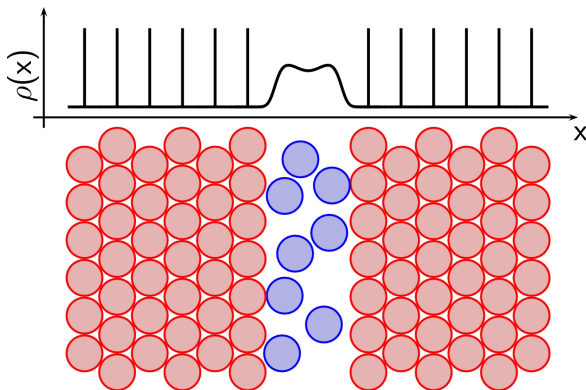
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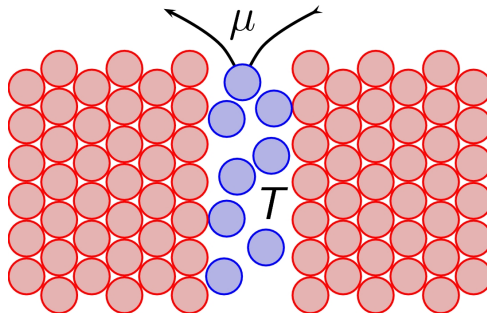
Observables and Parameters



- Key observables

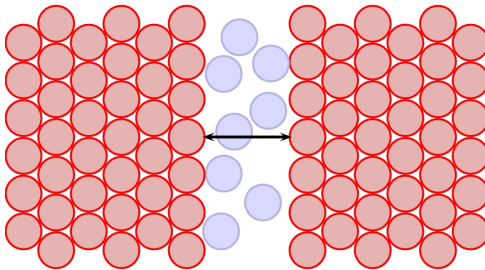
- ▶ Density profile $\rho(\vec{x})$
- ▶ Thermodynamic properties (grand potential, steric forces, pressure...)

Observables and Parameters



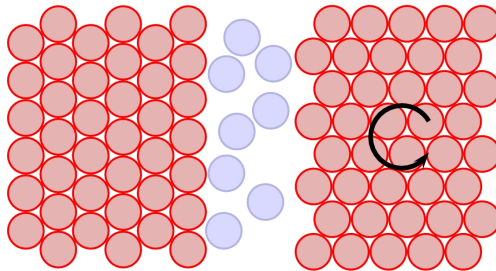
- Key observables: Density profile, potentials, pressure ...
- Parameters
 - ▶ temperature, chemical potential, ...
 - ▶ relative lattice orientation: gap, tilt, twist

Observables and Parameters



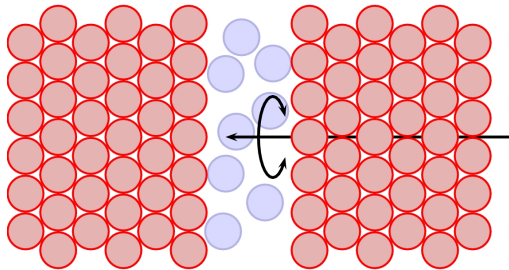
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The principles of Classical Density Functional Theory

Classical Density Functional Theory

- The **dimensionless, local potential** $u(\vec{x}) = \beta(\mu - U(\vec{x}))$ is a **unique functional** of the density $\rho(\vec{x})$.
- **Minimise the functional**

$$\widetilde{W}([u], [\tilde{\rho}]) = F([\tilde{\rho}]) - \int d^d x u(\vec{x}) \tilde{\rho}(\vec{x})$$

- Find $\tilde{\rho}_0$ that satisfies

$$\left. \frac{\delta}{\delta \tilde{\rho}} \right|_{\tilde{\rho} \equiv \tilde{\rho}_0} \widetilde{W}([u], [\tilde{\rho}]) \equiv 0$$

- At minimum: $\widetilde{W}([u], [\tilde{\rho} = \tilde{\rho}_0]) = W([u])$, the **grand potential**

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The rôle of the effective potential

Rôle of the effective potential $C(\vec{x})$

- **Effective potential** $C(\vec{x})$: Difference between ideal potential $\ln(\Lambda\rho(\vec{x}))$ and local potential $u(\vec{x})$
- **Standard approximation**:
Expand $C(\vec{x})$ in functional Taylor series
- **Perturbation theory** about the infinite, uniform liquid
- All interaction within the liquid enters solely through $C(\vec{x})$
→ **Direct correlation function**
- Test beds of increasing complexity:
Hard rods, hard spheres, Silicon

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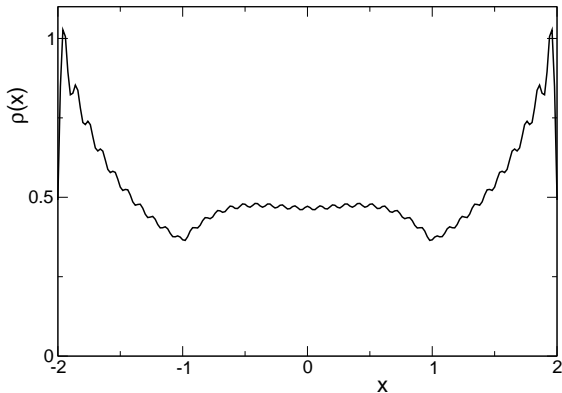
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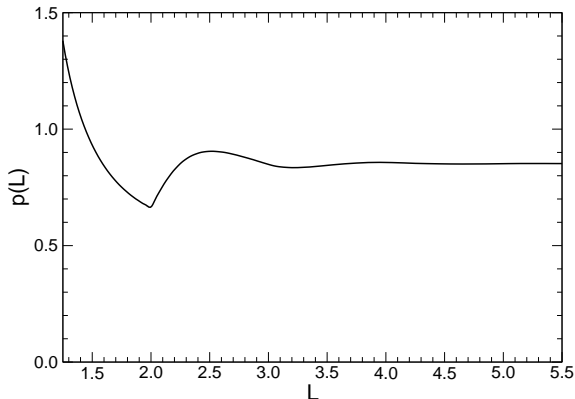
with S. Dungey and N. Gnanathas



Density profile for $L = 4$, $\rho_0 = 0.46$

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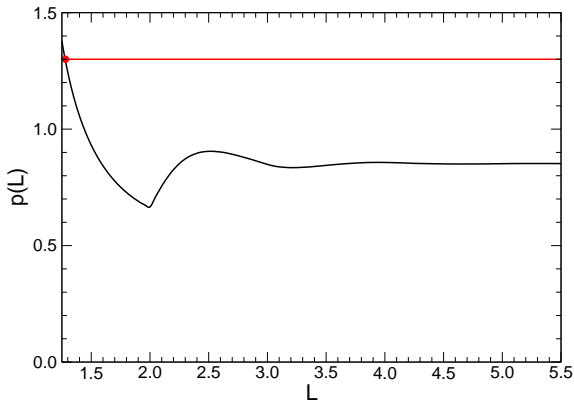
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Pressure as a function of wall separation

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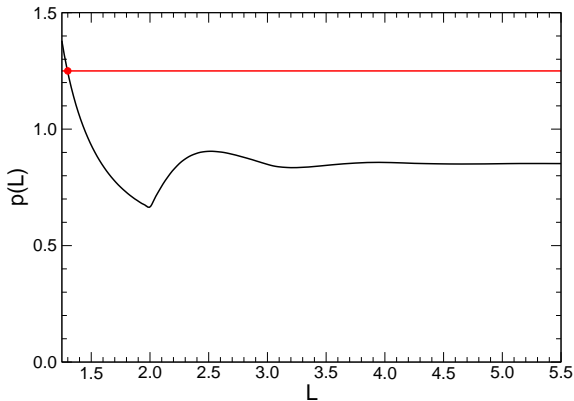
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Equilibrium width due to external walls

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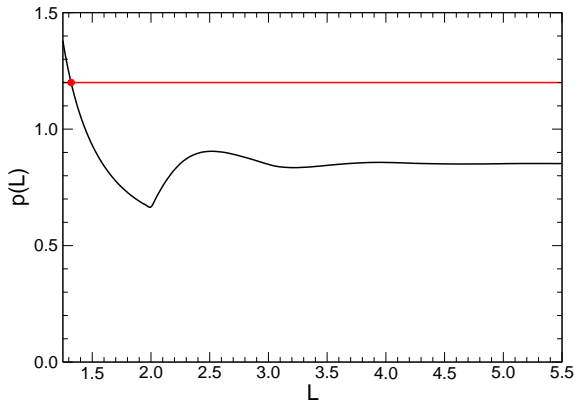
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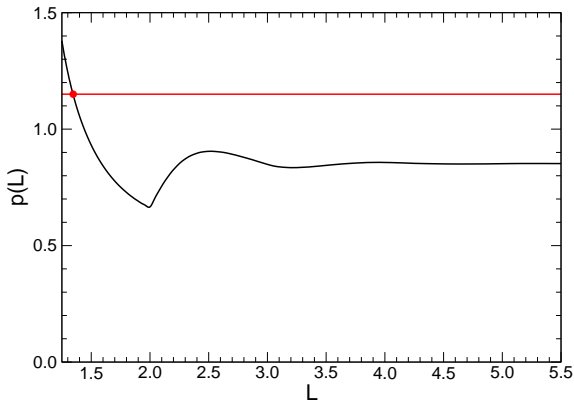
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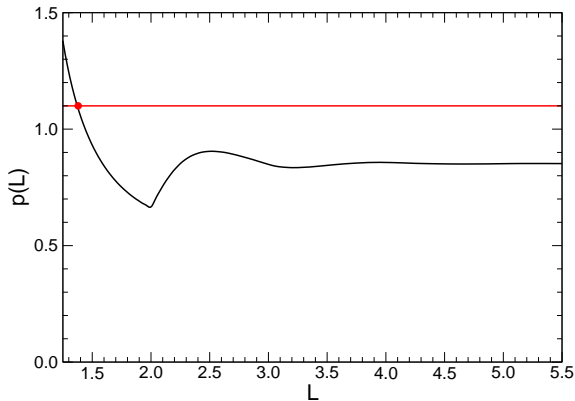
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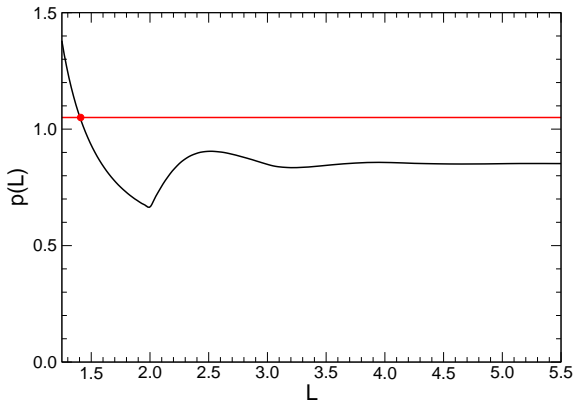
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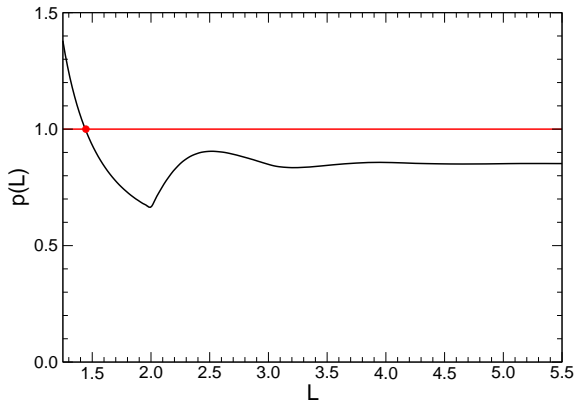
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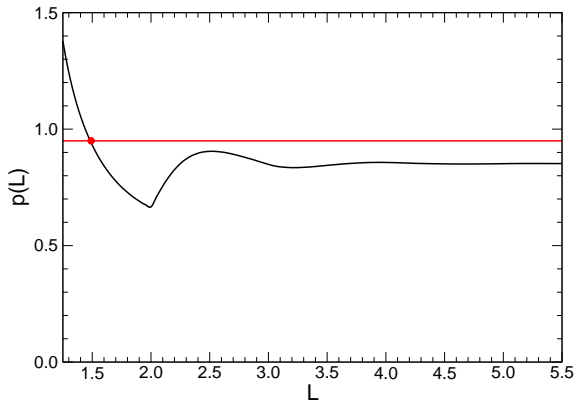
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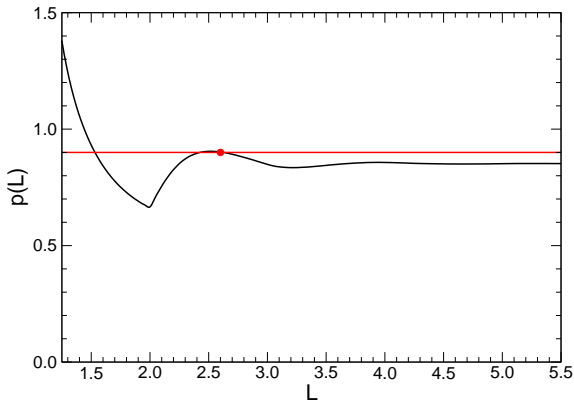
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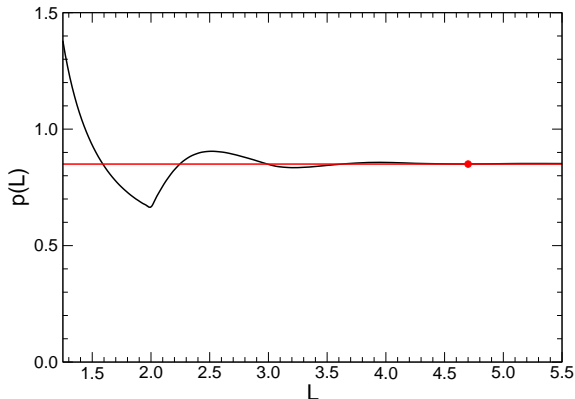
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Phase Field Models

Features of Phase Field Models

- Study dynamical and equilibrium properties of phase boundaries
- Avoid boundary problems and no-overhang approximation
- Phase field $\phi(\vec{x})$ describes “degree of liquid-ness”
- Provides equation of motion for phase changes and their interfaces
- Problem: Functional *ad hoc*, physical meaning of ϕ unclear

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Phase Field Model from DFT

Based on (Haymet and Oxtoby 1981, Oxtoby and Haymet 1982)

- **Standard method:** Write density ρ as a Fourier sum,

$$\rho(\vec{x}) = \rho_0 \left(1 + \eta + \sum_{\{\vec{k}\}} \mu_k e^{i\vec{k}\vec{x}} \right)$$

- Characterise liquid/solid by **spatially varying coefficients** $\mu_{\vec{k}}(\vec{x})$
- Oxtoby and Haymet:

Find solution of **integro-differential equation**

- Alternatively, **minimise** with respect to $\tilde{\eta}$ and $\tilde{\mu}_k$

$$\tilde{W} = \int d^d x f(\tilde{\eta}, \tilde{\mu}_n) + \frac{1}{4} c_0'' (\vec{\nabla} \tilde{\eta})^2 + \frac{1}{4} \sum_{\{\vec{k}\}} c_k'' (\hat{k} \vec{\nabla} \tilde{\mu}_k)^2 - u \tilde{\rho}$$

- $\tilde{\mu}_k(\vec{x})$ **is a phase field!** — Resulting **equation of motion**

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A. D. J. Haymet and David W. Oxtoby

A molecular theory of the solid-liquid interface

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A molecular theory of the solid-liquid interface. II. Study of bcc crystal-melt interfaces

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